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Classification
of Partial
Differential ...

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Introduction

Mathematics-II (Differential Equations) Lecture Notes April 16, 2020

by

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Example

Discuss D'Alembert's solution of one dimensional wave equation. or

Show that the general solution of the wave equation

$$c^2 \frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial t^2} \text{ is } u(x, t) = \phi(x + ct) + \psi(x - ct),$$

where ϕ and ψ are arbitrary functions.



Solution: Given equation is

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$$

Let v and w be two new independent variables such that

$$w = x + ct \quad \text{and} \quad v = x - ct \quad (1)$$

Now

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial w} \frac{\partial w}{\partial x} + \frac{\partial u}{\partial v} \frac{\partial v}{\partial x}$$



Using equation (1), we have

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial w} + \frac{\partial u}{\partial v} \quad \text{So that} \quad \frac{\partial}{\partial x} = \frac{\partial}{\partial w} + \frac{\partial}{\partial v} \quad (2)$$

Thus

$$\begin{aligned} \frac{\partial^2 u}{\partial x^2} &= \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) \implies \frac{\partial^2 u}{\partial x^2} = \left(\frac{\partial}{\partial w} + \frac{\partial}{\partial v} \right) \left(\frac{\partial u}{\partial w} + \frac{\partial u}{\partial v} \right) \\ &= \frac{\partial^2 u}{\partial w^2} + 2 \frac{\partial^2 u}{\partial w \partial v} + \frac{\partial^2 u}{\partial v^2} \end{aligned} \quad (3)$$

Again

$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial w} \frac{\partial w}{\partial t} + \frac{\partial u}{\partial v} \frac{\partial v}{\partial t}$$



Using equation (1), we have

$$\frac{\partial u}{\partial t} = c \frac{\partial u}{\partial w} - c \frac{\partial u}{\partial v} \quad \text{So that} \quad \frac{\partial}{\partial t} = c \left(\frac{\partial}{\partial w} - \frac{\partial}{\partial v} \right) \quad (4)$$

Thus

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial t} \left(\frac{\partial u}{\partial t} \right) \implies \frac{\partial^2 u}{\partial t^2} = c^2 \left(\frac{\partial}{\partial w} - \frac{\partial}{\partial v} \right) \left(\frac{\partial u}{\partial w} - \frac{\partial u}{\partial v} \right)$$

$$\begin{aligned} \frac{\partial^2 u}{\partial t^2} &= c^2 \left(\frac{\partial^2 u}{\partial w^2} - 2 \frac{\partial^2 u}{\partial w \partial v} + \frac{\partial^2 u}{\partial v^2} \right) \\ \implies \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} &= \left(\frac{\partial^2 u}{\partial w^2} - 2 \frac{\partial^2 u}{\partial w \partial v} + \frac{\partial^2 u}{\partial v^2} \right) \end{aligned} \quad (5)$$



Using (3) and (5) reduces to

$$\frac{\partial^2 u}{\partial w^2} + 2 \frac{\partial^2 u}{\partial w \partial v} + \frac{\partial^2 u}{\partial v^2} = \frac{\partial^2 u}{\partial w^2} - 2 \frac{\partial^2 u}{\partial w \partial v} + \frac{\partial^2 u}{\partial v^2} \implies \frac{\partial^2 u}{\partial w \partial v} = 0 \quad (6)$$

$$\frac{\partial}{\partial v} \left(\frac{\partial u}{\partial w} \right) = 0 \quad (7)$$



Integrating (7) w.r.t. v , we get

$$\frac{\partial u}{\partial w} = F(w), \quad (8)$$

where F is an arbitrary function of w .

Integrating (8) w.r.t. w , we get

$$u = \int F(w)dw + \psi(v),$$

where ψ is an function of v . Then

$$u = \phi(w) + \psi(v), \text{ where } \phi(w) = \int F(w)dw$$

or

$$u = \phi(x + ct) + \psi(x - ct).$$



General solution of one-dimensional heat (diffusion) equation satisfying the given boundary and initial conditions

Consider one-dimensional heat equation

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{k} \frac{\partial u}{\partial t},$$

where $u(x, t)$ is the temperature of the bar. If both the ends of a bar of length a are at temperature zero and initial temperature is to be prescribed function $f(x)$ in the bar, then find the temperature at a subsequent time t . More precisely, the faces $x = 0$ and $x = a$ of an infinite slab are maintained at zero temperature. Given that the temperature $u(x, t) = f(x)$ at $t = 0$. Find the temperature at a subsequent time t .



Solution: Given that

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{k} \frac{\partial u}{\partial t}, \quad (9)$$

with boundary conditions $u(0, t) = 0$, $u(a, t) = 0$.

The initial condition is given by $u(x, 0) = f(x)$, $0 < x < a$

Let the given equation has the solution of the form

$u(x, t) = X(x)T(t)$, where X is function of x alone and T is

function of t alone. Now $\frac{\partial^2 u}{\partial x^2} = X''(x)T(t)$ and

$\frac{\partial u}{\partial t} = X(x)T'(t)$. Putting these values in given equation, we have

$$X''T = \frac{1}{k}XT' \implies \frac{X''}{X} = \frac{T'}{kT}, \quad (10)$$



Since x and t are independent variables, therefore above equation can only true if each side is equal to the same constant. i.e.

$$\frac{X''}{X} = \frac{T'}{kT} = \mu(\text{constant}) \implies X'' - \mu X = 0 \text{ and} \\ T' - \mu kT = 0$$

These are ordinary differential equation of second order and first order with constant coefficient. Now to solve these two equations

$$X'' - \mu X = 0 \quad (11)$$

and

$$T' - \mu kT = 0. \quad (12)$$

Now three cases arises:



Case-I When $\mu = 0$, then both equations reduces to

$$X'' = 0 \implies X = a_1x + a_2$$

Using boundary conditions $u(0, t) = 0 = u(a, t)$, the trial solution $u(x, t) = X(x)T(t)$ becomes

$$0 = X(0)T(t) \quad \text{and} \quad 0 = X(a)T(t).$$

Since $T(t) = 0 \implies u(x, t) = 0$, so we suppose that $T(t) \neq 0$. Then we have $X(0) = 0$ and $X(a) = 0$. Now using these boundary conditions, the solution $X = a_1x + a_2$ becomes $0 = a_1 \cdot 0 + a_2$ and $0 = a_1 \cdot a + a_2 \implies a_1 = 0 = a_2$, so that $X(x) = 0$, which yields $u(x, t) = 0$. So we reject case-I, when $\mu = 0$.



Case-II When $\mu > 0$, we can take $\mu = \lambda^2$ (say), then equations $X'' - \mu X = 0$ reduces to

$$X'' - \lambda^2 X = 0 \implies \text{the auxiliary equation is}$$
$$(m^2 - \lambda^2) = 0 \implies m = \pm\lambda. \text{ Therefore its solution will be}$$
$$X = b_1 e^{\lambda x} + b_2 e^{-\lambda x}$$

Using boundary conditions $u(0, t) = 0 = u(a, t)$, the trial solution $u(x, t)X(x)T(t)$ becomes

$$0 = X(0)T(t) \quad \text{and} \quad 0 = X(a)T(t).$$

Since $T(t) = 0 \implies u(x, t) = 0$, so we suppose that $T(t) \neq 0$. Then we have $X(0) = 0$ and $X(a) = 0$. Now using these boundary conditions, the solution $X = b_1 e^{\lambda x} + b_2 e^{-\lambda x}$ becomes $0 = b_1 e^{\lambda \cdot 0} + b_2 e^{-\lambda \cdot 0}$ and $0 = b_1 e^{\lambda a} + b_2 e^{-\lambda a} \implies 0 = b_1 + b_2$ and $b_1 e^{\lambda a} + b_2 e^{-\lambda a} \implies b_1 = b_2 = 0$, so that $X(x) = 0$, which yields $u(x, t) = 0$. So again we reject case-II, when $\mu > 0$.



Case-III When $\mu < 0$, we can take $\mu = -\lambda^2$ (say), then first equation reduces to

$$X'' + \lambda^2 X = 0 \implies \text{the auxiliary equation is}$$
$$(m^2 + \lambda^2) = 0 \implies m = \pm \lambda i. \text{ Therefore its solution will be}$$
$$X = c_1 \cos(\lambda x) + c_2 \sin(\lambda x)$$

Using boundary conditions $u(0, t) = 0 = u(a, t)$, the trial solution becomes

$$0 = X(0)T(t) \quad \text{and} \quad 0 = X(a)T(t).$$

Since $T(t) = 0 \implies u(x, t) = 0$, so we suppose that $T(t) \neq 0$. Then we have $X(0) = 0$ and $X(a) = 0$. Now using these boundary conditions, the solution $X = c_1 \cos(\lambda x) + c_2 \sin(\lambda x)$ becomes $0 = c_1 \cos(\lambda \cdot 0) + c_2 \sin(\lambda \cdot 0)$ and $0 = c_1 \cos(\lambda a) + c_2 \sin(\lambda a) \implies c_1 = 0$ and $c_2 \sin(\lambda a) = 0$



Now for non-trivial solution of given wave equation, we can not take $c_2 = 0$

$$\implies \sin \lambda a = 0 \implies \lambda a = n\pi \quad n = 1, 2, 3, \dots$$

Thus $\lambda = \frac{n\pi}{a}, n = 1, 2, 3, \dots$

Hence non-zero solution $X_n(x)$ are given by

$$X_n(x) = (c_2)_n \sin\left(\frac{n\pi x}{a}\right) \quad (13)$$



Now the solution corresponding to the equation $T' + \lambda^2 k T = 0$ is

$$\frac{T'}{T} = -\lambda^2 k \quad (14)$$

By integrating we get

$$\log T = -\lambda^2 k t + \log c_3 \implies T = c_3 e^{-\lambda^2 k t} \implies T = c_3 e^{-(n^2 \pi^2 / a^2) t} \quad (15)$$

Hence solution is $T_n(t) = D_n e^{-C_n^2 t}$, where $C_n = (n^2 \pi^2 k / a^2)$ and $D_n = c_3$ are new arbitrary constants.

The general solution is

$$u_n(x, t) = \sum_{n=1}^{\infty} E_n \sin\left(\frac{n\pi x}{a}\right) e^{-C_n^2 t}, \quad (16)$$

where $E_n = (c_2)_n D_n$ is another new arbitrary constants.



Substituting $t = 0$ in (16) and using initial condition $u(x, 0) = f(x)$, we get

$$f(x) = \sum_{n=1}^{\infty} E_n \sin\left(\frac{n\pi x}{a}\right) \quad (17)$$

Which are Fourier sin series of expansion $f(x)$. Accordingly we get

$$E_n = \frac{2}{a} \int_0^a f(x) \sin \frac{n\pi x}{a} dx \quad (18)$$

Hence the required solution is given by the equation (16) and E_n given by the equation (18).



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Thanks !!!