

## <span id="page-0-0"></span>Lok Nayak Jai Prakash Institute of Technology Chapra, Bihar-841302

Classification of Partial [Differential ...](#page-25-0)

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Mathematics-II (Differential Equations) Lecture Notes April 15, 2020

by

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 $\mathcal{A} \cap \mathcal{A} \cap \mathcal{B} \longrightarrow \mathcal{A} \cap \mathcal{B} \longrightarrow \mathcal{B}$ 

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### [Introduction](#page-1-0) General solution of one-dimensional wave (vibrational) equation satisfying the given boundary conditions

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**[Introduction](#page-1-0)** 

Consider one-dimensional wave equation

$$
\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2},
$$

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with boundary conditions  $u(0, t) = 0$  and  $u(a, t) = 0$ ,  $\forall t$ .



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#### Solution: Given that

$$
\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2},\tag{1}
$$

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with boundary conditions  $u(0, t) = 0$  and  $u(a, t) = 0$ . Let the given equation has the solution of the form  $u(x,t) = X(x)T(t)$ , where X is function of x alone and T is function of t alone. Now  $\frac{\partial^2 u}{\partial x^2}$  $\frac{\partial^2 u}{\partial x^2} = X''(x)T(t)$  and  $\partial^2 u$  $\frac{\partial^2 u}{\partial t^2} = X(x)T''(t).$ 



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Putting these values in given equation, we have

$$
X''T = \frac{1}{c^2}XT'' \implies \frac{X''}{X} = \frac{T''}{c^2T},\tag{2}
$$

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Since  $x$  and  $t$  are independent variables, therefore above equation can only true if each side is equal to the same constant. i.e.

$$
\frac{X''}{X} = \frac{T''}{c^2T} = k(constant) \implies X'' - kX = 0 \text{ and}
$$

$$
T'' - c^2 kT = 0
$$

These are ordinary differential equation of second order with constant coefficient. Now to solve these two equations  $X'' - kX = 0$  and  $T'' - c^2 kT = 0$ , three cases arises:



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Case-1 When 
$$
k = 0
$$
, then both equations reduces to  

$$
X'' = 0 \implies X = a_1 x + a_2
$$

$$
T'' = 0 \implies T = a_3 t + a_4.
$$

Thus the required solution is

<span id="page-5-0"></span>
$$
u(x,t) = (a_1x + a_2)(a_3t + a_4). \tag{3}
$$

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**Case-II** When  $k > 0$ , we can take  $k = \lambda^2 (say)$ , then both equations reduces to

 $X''-\lambda^2 X=0\implies$  the auxiliary equation is  $(m^2 - \lambda^2) = 0 \implies m = \pm \lambda$ . Therefore its solution will be  $X = b_1 e^{\lambda x} + b_2 e^{-\lambda x}$ 

and

$$
T'' - c^2 \lambda^2 T = 0 \implies T = b_3 e^{c\lambda t} + b_4 e^{-c\lambda t}.
$$

Thus the required solution is

<span id="page-6-0"></span>
$$
u(x,t) = (b_1 e^{\lambda x} + b_2 e^{-\lambda x})(b_3 e^{c\lambda t} + b_4 e^{-c\lambda t}).
$$
 (4)

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**Case-III** When  $k < 0$ , we can take  $k = -\lambda^2 (say)$ , then both equations reduces to

 $X'' + \lambda^2 X = 0 \implies$  the auxiliary equation is  $(m^2 + \lambda^2) = 0 \implies m = \pm \lambda i$ . Therefore its solution will be  $X = c_1 \cos(\lambda x) + c_2 \sin(\lambda x)$ 

and

$$
T'' + c^2 \lambda^2 T = 0 \implies T = c_3 \cos(c\lambda t) + c_4 \sin(c\lambda t).
$$

Thus the required solution is

<span id="page-7-0"></span> $u(x,t) = (c_1 \cos(\lambda x) + c_2 \sin(\lambda x))(c_3 \cos(c\lambda t) + c_4 \sin(c\lambda t)).$ (5)

Thus the equation [\(3\)](#page-5-0), [\(4\)](#page-6-0) and [\(5\)](#page-7-0) are various possible solution of the given wave equation.  $2Q$ 



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Given boundary conditions are  $u(0,t) = u(a,t) = 0 \quad \forall t$  In view of the boundary condition, the solution given by the equation [\(3\)](#page-5-0) becomes

$$
0 = a_2(a_3t + a_4) \text{ and } 0 = (a_1a + a_2)(a_3t + a_2)
$$
  

$$
\implies a_2 = 0 \text{ and } (a_1a + a_2) = 0 \implies a_1 = a_2 = 0
$$

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Hence  $u(x,t) = 0 \quad \forall t$ . This is a trivial solution.



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Again, in view of the boundary condition, the solution given by the equation [\(4\)](#page-6-0) becomes

$$
0 = (b_1 + b_2)(b_3e^{c\lambda t} + b_4e^{-c\lambda t})
$$
 and  
\n
$$
0 = (b_1e^{\lambda a} + b_2e^{-\lambda a})(b_3e^{c\lambda t} + b_4e^{-c\lambda t})
$$
  
\n
$$
\implies (b_1 + b_2) = 0
$$
 and 
$$
b_1e^{\lambda a} + b_2e^{-\lambda a} = 0
$$
  
\n
$$
\implies b_1 = b_2 = 0
$$

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Hence  $u(x, t) = 0$   $\forall t$ . This is also a trivial solution.



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Finally, in view of the boundary condition, the solution given by the equation [\(5\)](#page-7-0) becomes

$$
0 = c_1(c_3 \cos(c\lambda t) + c_4 \sin(c\lambda t)) \text{ and}
$$
  

$$
0 = (c_1 \cos(\lambda a) + c_2 \sin(\lambda a))(c_3 \cos(c\lambda t) + c_4 \sin(c\lambda t))
$$

$$
\implies c_1 = 0 \quad \text{and} \quad c_2 \sin \lambda a = 0
$$

Now for non-trivial solution of given wave equation, we can not take  $c_2 = 0$ 

$$
\implies
$$
 sin  $\lambda a = 0 \implies \lambda a = n\pi$   $n = 1, 2, 3, ...$ 

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Thus 
$$
\lambda = \frac{n\pi}{a}
$$
,  $n = 1, 2, 3, ...$ 



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Hence the solution given by the equation [\(5\)](#page-7-0) becomes

$$
u_n(x,t) = c_2 \sin \frac{n\pi}{a} \left( c_3 \cos \frac{n\pi ct}{a} + c_4 \sin \frac{n\pi ct}{a} \right)
$$

$$
n = 1, 2, 3, ...
$$

$$
u_n(x,t) = \sin\frac{n\pi}{a} \left( E_n \cos\frac{n\pi ct}{a} + F_n \sin\frac{n\pi ct}{a} \right)
$$

$$
n = 1, 2, 3, \dots
$$

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Where  $E_n = (c_2c_3)$  and  $F_n = (c_2c_4)$  are new arbitrary constants.



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Since the given wave equation is linear, its most general solution is obtained by applying the principle of superposition, the required solution is

$$
u(x,t) = \sum_{n=1}^{\infty} u_n(x,t) =
$$
  

$$
\sum_{n=1}^{\infty} \sin \frac{n\pi}{a} \left( E_n \cos \frac{n\pi ct}{a} + F_n \sin \frac{n\pi ct}{a} \right) \quad n = 1, 2, 3, ...
$$

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[Introduction](#page-1-0)

### General solution of one-dimensional wave (vibrational) equation satisfying the given boundary and initial conditions

Consider one-dimensional wave equation

$$
\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2},
$$

where  $u(x, t)$  is the deflection of the string. the solution of this equation shows how the string moves. More precisely, if the ends of string are fixed at  $x = 0$  and  $x = a$ , we have the two boundary conditions.

$$
u(0,t)=0 \text{ and } u(a,t)=0, \quad \forall t.
$$

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The form of the motion of the string will depend on the initial deflection (deflection at  $t = 0$ ) and on the intial velocity (velocity at  $t = 0$ ). Denoting the initial deflection by  $f(x)$  and initial velocity by  $q(x)$ , we get two initial conditions

$$
u(x,0) = f(x), \qquad 0 \le x \le a
$$
  
and  

$$
\left(\frac{\partial u}{\partial t}\right)_{t=0} = g(x), \quad i.e. \quad u_t(x,0) = g(x) \qquad 0 \le x \le a.
$$

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Solution: Given that

$$
\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2},\tag{6}
$$

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with boundary conditions  $u(0, t) = 0$ ,  $u(a, t) = 0, u(x, 0) = f(x)$  and  $u_t(x, 0) = g(x), 0 \le x \le a$ . Let the given equation has the solution of the form  $u(x,t) = X(x)T(t)$ , where X is function of x alone and T is function of t alone. Now  $\frac{\partial^2 u}{\partial x^2}$  $\frac{\partial^2 u}{\partial x^2} = X''(x)T(t)$  and  $\partial^2 u$  $\frac{\partial^2 u}{\partial t^2} = X(x)T''(t)$ . Putting these values in given equation, we have  $X''T = \frac{1}{\epsilon}$  $\frac{1}{c^2}XT'' \implies \frac{X''}{X}$  $\frac{X''}{X} = \frac{T''}{c^2T}$  $c^2T$  $(7)$ 



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Since  $x$  and  $t$  are independent variables, therefore above equation can only true if each side is equal to the same constant. i.e.

$$
\frac{X''}{X} = \frac{T''}{c^2T} = k(constant) \implies X'' - kX = 0 \text{ and}
$$

$$
T'' - c^2kT = 0
$$

These are ordinary differential equation of second order with constant coefficient. Now to solve these two equations  $X'' - kX = 0$  and  $T'' - c^2 kT = 0$ , three cases arises:

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**Case-I** When  $k = 0$ , then both equations reduces to

$$
X'' = 0 \implies X = a_1 x + a_2
$$

Using boundary conditions  $u(0, t) = 0 = u(a, t)$ , the trial solution becomes

$$
0 = X(0)T(t) \qquad \text{and} \qquad 0 = X(a)T(t).
$$

Since  $T(t) = 0 \implies u(x, t) = 0$ , so we suppose that  $T(t) \neq 0$ . Then we have  $X(0) = 0$  and  $X(a) = 0$ . Now using these boundary conditions, the solution  $X = a_1x + a_2$  becomes  $0 = a_1.0 + a_2$  and  $0 = a_1.a + a_2 \implies a_1 = 0 = a_2$ , so that  $X(x) = 0$ , which yields  $u(x, t) = 0$ . So we reject case-I, when  $k=0.$ 

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<span id="page-18-0"></span>**Case-II** When  $k > 0$ , we can take  $k = \lambda^2 (say)$ , then first equations reduces to

 $X''-\lambda^2 X=0\implies$  the auxiliary equation is  $(m^2 - \lambda^2) = 0 \implies m = \pm \lambda$ . Therefore its solution will be  $X = b_1 e^{\lambda x} + b_2 e^{-\lambda x}$ 

Using boundary conditions  $u(0, t) = 0 = u(a, t)$ , the trial solution becomes

 $0 = X(0)T(t)$  and  $0 = X(a)T(t)$ .

Since  $T(t) = 0 \implies u(x, t) = 0$ , so we suppose that  $T(t) \neq 0$ . Then we have  $X(0) = 0$  and  $X(a) = 0$ . Now using these boundary conditions, the solution  $X=b_1e^{\lambda x}+b_2e^{-\lambda x}$  becomes  $0=b_1e^{\lambda.0}+b_2e^{-\lambda.0}$  and  $0=b_1e^{\lambda a}+b_2e^{-\lambda a} \implies 0=b_1+b_2$ and  $b_1e^{\lambda a}+b_2e^{-\lambda a} \implies b_1=b_2=0$ , so that  $X(x)=0$ , which yi[e](#page-1-0)lds $u(x,t) = 0$  $u(x,t) = 0$ . So again we reject [ca](#page-17-0)[se](#page-19-0)[-I](#page-17-0)[I,](#page-18-0) [w](#page-19-0)[h](#page-0-0)e[n](#page-25-0)  $k \ge 0$  $k \ge 0$ .

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**Case-III** When  $k < 0$ , we can take  $k = -\lambda^2 (say)$ , then first equations reduces to

 $X'' + \lambda^2 X = 0 \implies$  the auxiliary equation is  $(m^2 + \lambda^2) = 0 \implies m = \pm \lambda i$ . Therefore its solution will be  $X = c_1 \cos(\lambda x) + c_2 \sin(\lambda x)$ 

Using boundary conditions  $u(0, t) = 0 = u(a, t)$ , the trial solution becomes

 $0 = X(0)T(t)$  and  $0 = X(a)T(t)$ .

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Since  $T(t) = 0 \implies u(x, t) = 0$ , so we suppose that  $T(t) \neq 0$ . Then we have  $X(0) = 0$  and  $X(a) = 0$ . Now using these boundary conditions, the solution  $X = c_1 \cos(\lambda x) + c_2 \sin(\lambda x)$ becomes  $0 = c_1 \cos(\lambda.0) + c_2 \sin(\lambda.0)$  and  $0 = c_1 \cos(\lambda a) + c_2 \sin(\lambda a) \implies c_1 = 0$  and  $0 = c_2 \sin(\lambda a) = 0$ 

Now for non-trivial solution of given wave equation, we can not take  $c_2 = 0$ 

$$
\implies \sin \lambda a = 0 \implies \lambda a = n\pi \quad n = 1, 2, 3, \dots
$$

Thus  $\lambda = \frac{n\pi}{n}$  $\frac{a}{a}$ ,  $n = 1, 2, 3, ...$ Hence non-zero solution  $X_n(x)$  are given by

$$
(c_2)_n \sin\left(\frac{n\pi x}{a}\right) \tag{8}
$$

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Similarly the solution corresponding to the equation  $T'' + \lambda^2 T = 0$  is

$$
T_n(t) = (c_3)_n \cos \frac{n\pi ct}{a} + (c_4)_n \sin \frac{n\pi ct}{a} \tag{9}
$$

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Hence the required solution is

<span id="page-21-0"></span>
$$
u(x,t) = \sum_{n=1}^{\infty} \sin \frac{n\pi x}{a} \left( E_n \cos \frac{n\pi ct}{a} + F_n \sin \frac{n\pi ct}{a} \right) \tag{10}
$$

Where  $E_n = ((c_2)_n(c_3))$  and  $F_n = ((c_2)_n(c_4)_n)$  are new arbitrary constants.



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In order to find a solution which also satisfy  $u(x, 0) = f(x)$  and  $u_t(x, 0) = g(x)$ , We differentiate equation [\(10\)](#page-21-0) w.r.t. t,

<span id="page-22-0"></span>
$$
\frac{\partial u}{\partial t} = \sum_{n=1}^{\infty} \left\{ \sin \frac{n\pi x}{a} \left( \frac{-n\pi c}{a} E_n \sin \frac{n\pi ct}{a} + \frac{n\pi c}{a} F_n \cos \frac{n\pi ct}{a} \right) \right\}
$$
\n(11)  
\nPut  $t = 0$  in equation (10) and (11) and using initial equation  
\n $u(x, 0) = f(x)$  and  $u_t(x, 0) = g(x)$ , we get

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 $f(x) = \sum_{n=0}^{\infty}$  $n=1$  $E_n \sin \frac{n \pi x}{a}$ a (12)

$$
g(x) = \sum_{n=1}^{\infty} \frac{n\pi c F_n}{a} \sin \frac{n\pi x}{a}
$$
 (13)

Which are Fourier sin series of expansion  $f(x)$  and  $g(x)$ , respectively. Accordingly we get

<span id="page-23-0"></span>
$$
E_n = \frac{2}{a} \int_0^a f(x) \sin \frac{n \pi x}{a} dx \tag{14}
$$

and

and

<span id="page-23-1"></span>
$$
F_n = \frac{2}{n\pi c} \int_0^a g(x) \sin\frac{n\pi x}{a} dx
$$
 (15)



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Hence the required solution is given by the equation [\(10\)](#page-21-0) where  $E_n$  and  $F_n$  are given by the equation [\(14\)](#page-23-0) and [\(15\)](#page-23-1).

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